

On modified Newton iterations for SPARK methods applied to constrained systems in mechanics

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Abstract. The application of modified Newton iterations to the solution of SPARK methods applied to a large class of overdetermined differential-algebraic equations (ODAEs) is described in some details. These ODAEs include the formulation of systems in mechanics with holonomic and nonholonomic constraints.

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1. INTRODUCTION

We consider the following class of systems of implicit, partitioned, additive, and overdetermined differential-algebraic equations (ODAEs)

$$\frac{d}{dt}y = v(t, y, z), \quad \frac{d}{dt}p(t, y, z) = f(t, y, z, \psi) + r(t, y, \lambda), \quad 0 = g(t, y), \quad 0 = g_t(t, y) + g_y(y)v(t, y, z), \quad 0 = k(t, y, z) \quad (1)$$

where we assume that

$$\begin{pmatrix} p_z & -r_\lambda & -f_\psi \\ g_y v_z & 0 & 0 \\ k_z & 0 & 0 \end{pmatrix} \text{ is nonsingular,} \quad (2)$$

and that

$$\begin{pmatrix} p_z & -r_\lambda \\ g_y v_z & 0 \end{pmatrix} \text{ is nonsingular.} \quad (3)$$

The variable $t \in \mathbb{R}$ is the independent variable, $y \in \mathbb{R}^{n_y}$ and $z \in \mathbb{R}^{n_z}$ are the *differential* variables, $\lambda \in \mathbb{R}^{n_\lambda}$ and $\psi \in \mathbb{R}^{n_\psi}$ are the *algebraic* variables. The initial values (y_0, z_0) at t_0 are assumed to be given and consistent, i.e., the constraints in (1) must be satisfied. Sufficient differentiability of the functions v, p, f, r, g, k are also assumed to ensure existence and uniqueness of a solution. The ODAEs (1) include the formulation of mechanical systems with mixed constraints of holonomic, nonholonomic, scleronomic, and rheonomic types. In mechanics the quantities y, v, p represent respectively certain coordinates, their velocities, and their momenta; the right-hand side of the second equations in (1) contains forces acting on the system; the corresponding ODAEs can be derived from the Lagrange-d'Alembert principle; λ and ψ are Lagrange multipliers associated respectively to the holonomic constraints $0 = g(t, y)$, $0 = g_t(t, y) + g_y(y)v(t, y, z)$ and to the nonholonomic constraints $0 = k(t, y, z)$.

2. (s, s) -SPARK METHODS

One step of an (s, s) -SPARK method applied to the system of ODAEs (1) with consistent initial values (y_0, z_0) at t_0 and stepsize h is given as follows

$$Y_i = y_0 + h \sum_{j=1}^s a_{ij} v(T_j, Y_j, Z_j) \quad \text{for } i = 1, \dots, s,$$

$$\begin{aligned}
p(T_i, Y_i, Z_i) &= p(t_0, y_0, z_0) + h \sum_{j=1}^s \widehat{a}_{ij} f(T_j, Y_j, Z_j, \Psi_j) + h \sum_{j=0}^s \widetilde{a}_{ij} r(\widetilde{T}_j, \widetilde{Y}_j, \widetilde{\Lambda}_j) \quad \text{for } i = 1, \dots, s, \\
\widetilde{Y}_i &= y_0 + h \sum_{j=1}^s \widetilde{a}_{ij} v(T_j, Y_j, Z_j) \quad \text{for } i = 0, 1, \dots, s, \\
0 &= g(\widetilde{T}_i, \widetilde{Y}_i) \quad \text{for } i = 0, 1, \dots, s, \\
0 &= \sum_{j=1}^s b_j c_j^{i-1} k(T_j, Y_j, Z_j) \quad \text{for } i = 1, \dots, s-1, \\
y_1 &= y_0 + h \sum_{j=1}^s b_j v(T_j, Y_j, Z_j), \\
p(t_1, y_1, z_1) &= p(t_0, y_0, z_0) + h \sum_{j=1}^s \widehat{b}_j f(T_j, Y_j, Z_j, \Psi_j) + h \sum_{j=0}^s \widetilde{b}_j r(\widetilde{T}_j, \widetilde{Y}_j, \widetilde{\Lambda}_j), \\
0 &= g(t_1, y_1), \\
0 &= g_t(t_1, y_1) + g_y(t_1, y_1) v(t_1, y_1, z_1), \\
0 &= k(t_1, y_1, z_1)
\end{aligned}$$

where

$$t_1 := t_0 + h, \quad T_i := t_0 + c_i h \quad \text{for } i = 1, \dots, s, \quad \widetilde{T}_i := t_0 + \widetilde{c}_i h \quad \text{for } i = 0, 1, \dots, s.$$

For Lobatto coefficients a similar definition was proposed in [1, 2]. The definition given here is more general as it can also include for example Gauss and Radau coefficients. We have four sets of coefficients (b_j, a_{ij}, c_i) , $(\widehat{b}_j, \widehat{a}_{ij})$, $(\widetilde{b}_j, \widetilde{a}_{ij})$, $(\widetilde{a}_{ij}, \widetilde{c}_i)$, where we have defined

$$c_i := \sum_{j=1}^s a_{ij} \quad \text{for } i = 1, \dots, s, \quad \widetilde{c}_i := \sum_{j=1}^s \widetilde{a}_{ij} \quad \text{for } i = 0, 1, \dots, s.$$

We assume that $\widetilde{a}_{0j} = 0$ for $j = 1, \dots, s$ which implies that $\widetilde{Y}_0 = y_0$, $\widetilde{c}_0 = 0$, $\widetilde{T}_0 = t_0$, and $0 = g(\widetilde{T}_0, \widetilde{Y}_0) = g(t_0, y_0)$ is thus automatically satisfied. We also assume that $\widetilde{a}_{sj} = b_j$ for $j = 1, \dots, s$ which implies that $\widetilde{Y}_s = y_1$, $\widetilde{c}_s = 1$, and $\widetilde{T}_s = t_1$. Hence, from $0 = g(\widetilde{T}_s, \widetilde{Y}_s)$ the condition $0 = g(t_1, y_1)$ is also automatically satisfied. Notice that the coefficients $(b_j, c_j)_{j=1}^s$ and $(\widetilde{b}_j, \widetilde{c}_j)_{j=0}^s$ generally correspond to two distinct quadrature formulas. We assume $b_i \neq 0$, $c_i \neq c_j$ for $i \neq j$, and the matrix A to be invertible. We use the following notation, $\mathbb{1}_s := (1, 1, \dots, 1)^T \in \mathbb{R}^s$, $\mathbf{0}_s := (0, 0, \dots, 0)^T \in \mathbb{R}^s$, $e_{s+1} := (0, 0, \dots, 0, 1)^T \in \mathbb{R}^{s+1}$, $I_s := \text{diag}(1, 1, \dots, 1) \in \mathbb{R}^{s \times s}$, $C := \text{diag}(c_1, c_2, \dots, c_s) \in \mathbb{R}^{s \times s}$, and we define

$$\alpha := \begin{pmatrix} A \\ b^T \end{pmatrix} \in \mathbb{R}^{(s+1) \times s}, \quad \widehat{\alpha} := \begin{pmatrix} \widehat{A} \\ \widehat{b}^T \end{pmatrix} \in \mathbb{R}^{(s+1) \times s}, \quad \widetilde{\alpha} := \begin{pmatrix} \widetilde{A} \\ \widetilde{b}^T \end{pmatrix} \in \mathbb{R}^{(s+1) \times (s+1)}.$$

We assume that $\widetilde{\alpha}$ is invertible. We also define

$$\widetilde{Q} := \begin{pmatrix} I & \mathbf{0}_s \end{pmatrix} + A M^{-1} \mathbb{1}_s \begin{pmatrix} -b^T A^{-1} & 1 \end{pmatrix} \in \mathbb{R}^{s \times (s+1)}$$

where we assume that

$$M := \begin{pmatrix} b^T \\ b^T - b^T A \\ \vdots \\ b^T - (s-1)b^T C^{s-2} A \end{pmatrix} \in \mathbb{R}^{s \times s} \quad \text{is invertible.}$$

We define

$$\gamma^T = \begin{pmatrix} \widetilde{\gamma}^T & \gamma_{s+1} \end{pmatrix}^T := \gamma_{s+1} \begin{pmatrix} -b^T A^{-1} & 1 \end{pmatrix} \neq 0 \in \mathbb{R}^{s+1}$$

which satisfies the s orthogonality conditions $\gamma^T \alpha = 0$. We define the invertible matrix Q by

$$Q := \begin{pmatrix} \widetilde{Q} \\ \gamma^T \end{pmatrix} \in \mathbb{R}^{(s+1) \times (s+1)}$$

and the matrix \check{Q} by

$$\check{Q} := Q \begin{pmatrix} \check{A}^{-1} & 0_s \\ 0_s^T & 1 \end{pmatrix} \in \mathbb{R}^{(s+1) \times (s+1)}$$

where we assume that

$$\check{A} := \begin{pmatrix} \bar{a}_{11} & \cdots & \bar{a}_{1s} \\ \vdots & \ddots & \vdots \\ \bar{a}_{s1} & \cdots & \bar{a}_{ss} \end{pmatrix} \in \mathbb{R}^{s \times s} \quad \text{is invertible}$$

and that $e_{s+1}^T Q \hat{\alpha} = 0_s$, for example by having $\hat{\alpha} = \alpha$. More details can be found in [3, 4, 5].

2.1. Reformulation of (s, s) -SPARK methods

To solve the nonlinear system of equations for (s, s) -SPARK methods we consider the application of modified Newton methods. In order to obtain an efficient implementation requiring only the decomposition of the 2 matrices in (2) and (3) we reformulate the nonlinear system of equations of (s, s) -SPARK methods equivalently as follows

$$0 = \begin{pmatrix} Y_1 \\ \vdots \\ Y_s \\ y_1 \end{pmatrix} - \mathbb{1}_{s+1} \otimes y_0 - h(\alpha \otimes I_{n_y}) \begin{pmatrix} v(T_1, Y_1, Z_1) \\ \vdots \\ v(T_s, Y_s, Z_s) \end{pmatrix}, \quad (4)$$

$$0 = (Q \otimes I_{n_z}) \left(\begin{pmatrix} p(T_1, Y_1, Z_1) \\ \vdots \\ p(T_s, Y_s, Z_s) \\ p(t_1, y_1, z_1) \end{pmatrix} - \mathbb{1}_{s+1} \otimes p(t_0, y_0, z_0) - h(\hat{\alpha} \otimes I_{n_z}) \begin{pmatrix} f(T_1, Y_1, Z_1, \Psi_1) \\ \vdots \\ f(T_s, Y_s, Z_s, \Psi_s) \end{pmatrix} - h(\tilde{\alpha} \otimes I_{n_z}) \begin{pmatrix} r(\tilde{T}_1, \tilde{Y}_1, \tilde{\Lambda}_1) \\ \vdots \\ r(\tilde{T}_s, \tilde{Y}_s, \tilde{\Lambda}_s) \end{pmatrix} \right), \quad (5)$$

$$0 = (\check{Q} \otimes I_{n_\lambda}) \begin{pmatrix} \frac{1}{h} g(\tilde{T}_1, y_0 + h \sum_{j=1}^s \bar{a}_{1j} v(T_j, Y_j, Z_j)) \\ \vdots \\ \frac{1}{h} g(\tilde{T}_s, y_0 + h \sum_{j=1}^s \bar{a}_{sj} v(T_j, Y_j, Z_j)) \\ g_r(t_1, y_1) + g_y(t_1, y_1) v(t_1, y_1, z_1) \end{pmatrix}, \quad (6)$$

$$0 = (\tilde{Q} \otimes I_{n_\psi}) \begin{pmatrix} k(T_1, Y_1, Z_1) \\ \vdots \\ k(T_s, Y_s, Z_s) \\ k(t_1, y_1, z_1) \end{pmatrix} \quad (7)$$

where

$$\begin{pmatrix} \tilde{Y}_0 \\ \tilde{Y}_1 \\ \vdots \\ \tilde{Y}_s \end{pmatrix} := \mathbb{1}_{s+1} \otimes y_0 - h(\bar{A} \otimes I_{n_y}) \begin{pmatrix} v(T_1, Y_1, Z_1) \\ \vdots \\ v(T_s, Y_s, Z_s) \end{pmatrix}.$$

2.2. Modified Jacobian and modified Newton iterations

The modified Jacobian of the nonlinear system of equations (4)-(5)-(6)-(7) can be taken as

$$\begin{pmatrix} I_{s+1} \otimes I_{n_y} & O & O & O \\ O & Q \otimes p_z & -hQ\tilde{\alpha} \otimes r_\lambda & -hQ\tilde{\alpha} \otimes f_\psi \\ O & Q \otimes g_y v_z & O & O \\ O & \tilde{Q} \otimes k_z & O & O \end{pmatrix}$$

where the partial derivatives $p_z, f_\psi, r_\lambda, g_y, k_z, v_z$ are evaluated for example at $t_0, y_0, z_0, \lambda_0, \psi_0$. We consider intermediate quantities to solve the linear systems of the modified Newton iterations. First we can obtain a block diagonal linear system with s matrix blocks (2) of dimension $n_z + n_\lambda + n_\psi$ for

$$\begin{pmatrix} \Delta_1^z \\ \vdots \\ \Delta_s^z \end{pmatrix} := (\tilde{Q} \otimes I_{n_z}) \begin{pmatrix} \Delta Z_1 \\ \vdots \\ \Delta Z_s \\ \Delta z_1 \end{pmatrix}, \quad \begin{pmatrix} \Delta_1^\lambda \\ \vdots \\ \Delta_s^\lambda \end{pmatrix} := h(\tilde{Q}\tilde{\alpha} \otimes I_{n_\lambda}) \begin{pmatrix} \Delta\tilde{\Lambda}_0 \\ \Delta\tilde{\Lambda}_1 \\ \vdots \\ \Delta\tilde{\Lambda}_s \end{pmatrix}, \quad \begin{pmatrix} \Delta_1^\psi \\ \vdots \\ \Delta_s^\psi \end{pmatrix} := h(\tilde{Q}\tilde{\alpha} \otimes I_{n_\psi}) \begin{pmatrix} \Delta\Psi_1 \\ \vdots \\ \Delta\Psi_s \end{pmatrix}.$$

By invertibility of $\tilde{Q}\tilde{\alpha}$ we obtain the values $\Delta\Psi_1, \dots, \Delta\Psi_s$ from

$$\begin{pmatrix} \Delta\Psi_1 \\ \vdots \\ \Delta\Psi_s \end{pmatrix} = \frac{1}{h} ((\tilde{Q}\tilde{\alpha})^{-1} \otimes I_{n_\psi}) \begin{pmatrix} \Delta_1^\psi \\ \vdots \\ \Delta_s^\psi \end{pmatrix}.$$

We also obtain a linear system of dimension $n_z + n_\lambda$ with matrix (3) for

$$\Delta_{s+1}^z := (\gamma^T \otimes I_{n_z}) \begin{pmatrix} \Delta Z_1 \\ \vdots \\ \Delta Z_s \\ \Delta z_1 \end{pmatrix}, \quad \Delta_{s+1}^\lambda := h(\gamma^T \tilde{\alpha} \otimes I_{n_\lambda}) \begin{pmatrix} \Delta\tilde{\Lambda}_0 \\ \Delta\tilde{\Lambda}_1 \\ \vdots \\ \Delta\tilde{\Lambda}_s \end{pmatrix}.$$

By invertibility of Q and $\tilde{\alpha}$ we then obtain the values $\Delta Z_1, \dots, \Delta Z_s, \Delta z_1$ and $\Delta\tilde{\Lambda}_0, \Delta\tilde{\Lambda}_1, \dots, \Delta\tilde{\Lambda}_s$ from

$$\begin{pmatrix} \Delta Z_1 \\ \vdots \\ \Delta Z_s \\ \Delta z_1 \end{pmatrix} = (Q^{-1} \otimes I_{n_z}) \begin{pmatrix} \Delta_1^z \\ \vdots \\ \Delta_s^z \\ \Delta_{s+1}^z \end{pmatrix}, \quad \begin{pmatrix} \Delta\tilde{\Lambda}_0 \\ \Delta\tilde{\Lambda}_1 \\ \vdots \\ \Delta\tilde{\Lambda}_s \end{pmatrix} = \frac{1}{h} ((Q\tilde{\alpha})^{-1} \otimes I_{n_\lambda}) \begin{pmatrix} \Delta_1^\lambda \\ \vdots \\ \Delta_s^\lambda \\ \Delta_{s+1}^\lambda \end{pmatrix}.$$

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